Effect of Grain Size on Electric and Magnetic Properties of Barium Hexaferrite (BaFe₁₂O₁₉)

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Polycrystalline Barium hexaferrite (BaFe₁₂O₁₉) having variable grain sizes were prepared by auto-combustion and solid state technique. SEM micrograph shows the variation of grain size (from nano range to bulk) of the system prepared at different temperatures. For this work electric transport (in the frequency range 100Hz to 1Mz) and magnetic study (M-H loop) carried out at room temperature. From the electric study it was found that, resistivity of the system varies inversely with grain size, however saturation magnetization (M_s) varies directly with grain size. In our report we were attempted to resolve the relation between grain size with electric and magnetic properties of the studied system.

Keywords: Microstructure, Impedance, Modulus, Magnetization

1. Introduction:

Barium hexaferrite (BaFe12O19) is widely used as a permanent magnet material as it possesses large saturation magnetization (M s), high coercive force (Hc) and high magnetic anisotropy (Ha) along with excellent chemical stability and corrosion resistivity [1–4]. Apart from large scale applications in permanent magnets such as: motors, actuators, sensors, isolators, circulators, phase shif ters, and miniature antennas etc., its ultrafi ne particle is distinctly used for high density recording media to increase the storage capacity and reduce medium noise [5,6]. BaM system retains hexagonal structure consists of a close-packed stacking of oxygen layers, with the iron (Fe) atoms distributed within three kinds of octahedral sites (12k, 2a and 4f2), one tetrahedral site (4f1), and one bi-pyramidal site (2b) which are coupled by super exchange interactions through the O^{2-} ions [12]. The magnetic and electrical properties of the BaM system depend on the distribution of the cations in octahedral, tetrahedral and bi-pyramidal sites [12-21].

2. Experimental Techniques:

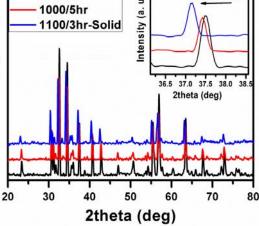
Barium hexaferrite nano range particles were prepared by the auto combustion method. Analytical grades Iron Nitrate Fe (NO₃).9H₂O and Barium Nitrate Ba (NO₃)₂ were taken as oxidants where Glycine $(C_2H_5NO_2)$ used as fuel for combustion. The stoichiometric amounts of these nitrates dissolved in distilled water and stirred at 150°C for proper homogenous mixture. After 2-3h stirring all the water evaporates and combustion taken place. Finally brown like powder was collected. The powder was calcined at 900°C for 12h and then formation of BaM pure phase was confirmed by the Rigaku Ultima-IV X-ray diffractometer (Cu target). For electric study the calcinated powder pressed into cylindrical pellets and sintered at 950°C for 5h then pure phase formation of BaM confirmed by X-ray diffractometer. For increasing of grain size we sinter the same pellet at 1000°C for 5h and 1100°C for 3h. The surface morphology of the BaM nanoparticle was estimated by the Nova Nano SEM-450 Field emission scanning electron microscopy (FESEM). Room temperature impedance with respect to frequency was measured by using HIOKI impedance analyzer (model IM3570) from 100 Hz to 1 MHz. The room temperature M-H loop study for all the systems done by SQUID VSM (Quantum Design).

3. Result and Discussion

950/5hr Peak shifting ntensity (a. u.) 1000/5hr 1100/3hr-Solid Intensity (a. u.) 2theta (deg)

3.1 Structural and Microstructural Study:

Fig.1



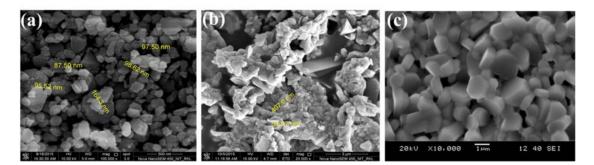


Fig. 2

From the Fig. 1 it was confirmed that all the X-ray diffraction peaks are matched with pure phase of BaM system without appearing of any extra peak. We can also observe that (insight of Fig. 1) peak shifts towards left with rise of temperature which is the indication of increasing of grain growth. From SEM micrograph it was well observed that the grains densely packed inside the system. It was also confirmed that grain size increases from nano range to micro rage as the sintering temperature increases. Here Fig. 2a, 2b and 2c sintered at 950°C, 1000°C and 1100°C respectively.

3.2 Impedance and Modulus Study:

From the Fig. it was observed that as grain size increases the resistive part (Z') of systems decreases. Here we can expect that, increasing of grain size can reduces the grain boundary density for a system which has higher restive nature than grain hence resistance of the system reduced. From the reactive nature (Z'') of the systems we can predict that all the systems have relaxation behaviour at room temperature. But quantitatively, the behaviour of Z' and Z'' with frequency are not exactly same. For examine the intrinsic and extrinsic contribution of the relaxation behaviour we taken the help of cole-cole of impedance and modulus spectroscopic study.

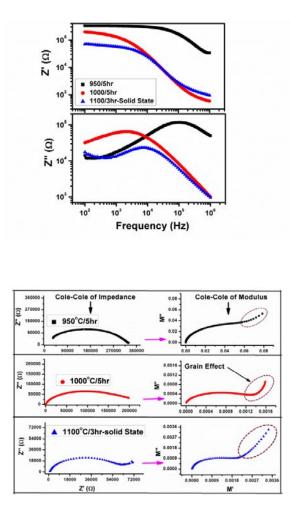


Fig. shows the cole-cole impedance and modulus at room temperature (RT). The cole-cole of impedance graph single relaxation behaviour due to which only single semicircle appears in the graph. Pattanayk et.al worked on bulk BaM system and found co-contribution of grain and grain boundary relaxation effect at RT which was confirmed from cole-cole of modulus study. Here if we observer the cole-cole of modulus spectrum of bulk BaM system (which is fired at 1100°C) then we then we can find similar type of behaviour (i.e. a semicircle connected with a larger tail). So the tail like feature (circular mark in the figure) in the graph is the intrinsic contribution (grain) where the semicircle is due to extrinsic (grain boundary) relaxation behaviour. But it is interesting that, when the grain size reduces (i.e. sample prepared at low temperature) the tail like feature (grain effect) gradually disappears. So here we predict that as grain size reduces the short range hopping conduction dominates over long range conduction inside the grain may due to increasing of hopping distance.

3.3 Magnetic Study:

Fig. 12 shows the hysteresis (M-H) loops of all systems at room temperature. It can be observe that the saturation magnetization (Ms), remnant magnetization (Mr) and coercive field (Hc) increases with rise of temperature (values are given in Table 4). The saturation magnetization of samples is derived from the law of saturation using the following equation (2) [i]:

$$M = M_s \left(1 - \frac{A}{H} - \frac{B}{H^2} \right) + \chi_p H$$
(2)

Where M_s is the saturation magnetization, A is inhomogeneity parameter, χ_P is the high field susceptibility and B is the anisotropy parameter. For hexagonal crystal structure, B can be expressed as equation (3):

$$B = \frac{H_a^2}{15} = \frac{4K_1^2}{15M_s^2} \dots (3)$$

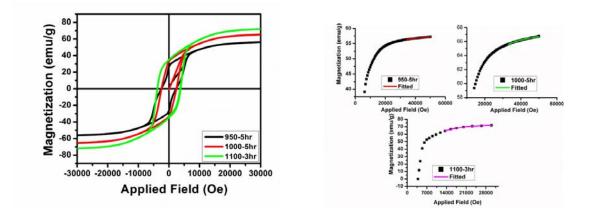
Where H_a is anisotropy field and K_1 is magneto-crystalline anisotropy constant (intrinsic magnetic parameter). The magneto-crystalline anisotropy constant can be calculate from quantitative analyses of the M-H curve by using the law of approach to saturation (LAS). The LAS can be written as (equation 4) [ii]:

$$M = M_s \left(1 - \frac{8K_1^2}{105\mu_0^2 M_s^2 H^2} \right) + cH$$
(4)

 M_s is the saturation magnetization, μ_0 is the magnetic permeability in free space, and c is linear constant term. The phenomenological linear term cH is very small at high magnetic field. The M-H loop was fitted (insights of Fig. 9) by using the equation (4) and magnetocrystalline anisotropy constants (K_1) were estimated (Table 3). Then anisotropy field (H_a) values were for both the system calculated by using equation 3.

It was observed that (Table 3) for composite system H_a was reduced as compared with BaM system. From the above discussion it can be predict that, presences of BaM-NBT interfaces in the system pinning of magnetization at the grain boundaries as the result anisotropy field

reduced. When the anisotropy field reduced it lowers the activation energy barrier by which domain walls will free to move hence coercivity of the composite system reduced [41].



Parameters \rightarrow	Ms	Mr	H _c (Oe)	Anisotropy	H _a (kOe)
Compositions↓	(emu/g)	(emu/g)		Constant (K_1) (erg/cm ³)	
950-5hr	57	29	2425	1.19E6	41.75
1000-5hr	67	33	2598	1.38E6	41.19
1100-3hr	73	34	3811	1.21E6	33.61

From Table- it can be observer that the saturation magnetization increases with grain size which may be due to larger domain size (decreasing of domain density) with less random orientation having less anisotropic field (H_a). But here we observe a contradictory result that with decreasing of Ha increasing of Hc. It was well known that if Ha reduces the domain wall strength will be reduces which can help to easy ration of domain along the field direction.

4. Conclusion:

Polycrystalline Barium hexaferrite (BaFe₁₂O₁₉) having variable grain sizes were prepared by using auto-combustion and solid state technique. SEM micrograph shows the variation of grain size (from nano range to bulk) of the system prepared at different temperatures. From the complex impedance spectroscopy it was found that grain relaxation evolves when the grain size of system increase. From magnetic study (M-H loop) it was found that as grain size increase the saturation magnetization increases.

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References:

[1] S. Castro, M.G. Rivas, J.M. Greneche, J. Mira, C. Rodrlguez, Journal of Magnetism and

Magnetic Material 152 (1996) 61-69.

[2] A. Mali, A. Ataie, Scripta Mater. 53 (2005) 1065.

[3] Vivek A. Rane, Sher Singh Meena, Suresh P. Gokhale, S.M. Yusuf, Girish J. Phatak, and Sadgopal. Date, Journal of Electronic Materials, 42 (2013) 4.

[4] Y.Y. Meng, M.H. He, Q. Zeng, D.L. Jiao, S. Shukla, R.V. Ramanujan, Z.W. Liu, Journal

of Alloys and Compounds 583 (2014) 220-225.

[5] T. Fujiwara, IEEE Transactions 21 (1985) 1480–1485.

[6] K. Yamamori, T. Suzuki, T.M. Fujiwara, IEEE Transactions 22 (1986) 1188–1190.

[12] F. M. M. Pereira, C. A. R. Junior, M. R. P. Santos, R. S. T. M. Sohn, F. N. A. Freire, J.

M. Sasaki, J. A. C. de Paiva, A. S. B. Sombra, J Mater Sci: Mater Electron, 19 (2008) 627–638.

[13] H. S ozeri, I. Kucuk, H. O zkan, Journal of Magnetism and Magnetic Materials 323(2011) 1799–1804.

[14] S.M. El-Sayed a,n, T.M. Meaz b,1, M.A. Amer b,2, H.A. El Shersaby, Physica B 426(2013) 137–143.

[15] Cong-Ju Li, BinWang, Jiao-NaWang, Journal of Magnetism and Magnetic Materials 324(2012) 1305–1311.

[16] Sonal Singhal, Tsering Namgyal, Jagdish Singh, Kailash Chandra, Sandeep Bansal, Ceramics International 37 (2011) 1833–1837. [17] Ashima, Sujata Sanghi, Ashish Agarwal, Reetu, Journal of Alloys and Compounds 513(2012) 436–444.

[18] Daming Chen, Yingli Liu, Yuanxun Li, Kai Yang, Huaiwu Zhang, Journal of Magnetism and Magnetic Materials 337–338 (2013) 65–69.

[19] A. Haq, M. Anis-ur-Rehman, Physica B 407 (2012) 822-826.

[20] Vaishali V. Soman, V.M. Nanoti, D.K. Kulkarni, Ceramics International 39 (2013)5713–5723.

[21] H. Sözeri, H. Deligöz, H. Kavas, A. Baykal, Ceramics International 40 (2014) 8645– 8657.

[22] Ashima, Sujata Sanghi, Ashish Agarwal, Reetu, Neetu Ahlawat, J. Appl. Phys. 112(2012) 014110.

[7] A.N.Bogdanov & I.E.Dragunov 1998

[[]i] Grössinger R. J Magn Magn Mater, 28 (1982) 137.

[[]ii] S. Chikazumi, Physics of Ferromagnetism, 2nd ed. (Oxford University Press, Oxford, 1997).

^[41] B.K.Bammannavar, L.R.Naik, J. Magn.Magn.Mater. 321 (2009) 382-387.